

Conductivity problems in random media

Nicolas CLOZEAU

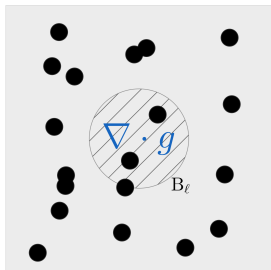


Joint work with Lihan WANG (CMU)

CANUM 2024

May 27 - May 31, 2024

Conductivity in random media



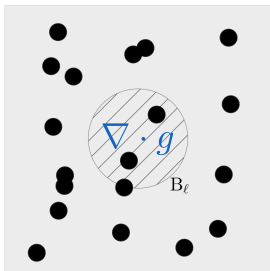
Goal: Compute the electric field generated by a localized charge in random heterogeneous media

Model: Compute ∇u where $u : \mathbb{R}^3 \rightarrow \mathbb{R}$ solves

$$-\nabla \cdot a \nabla u = \nabla \cdot g \quad \text{in } \mathbb{R}^3 \quad \text{such that } u(x) \xrightarrow{|x| \uparrow \infty} 0,$$

with $a = a(x)$ the (heterogeneous) conductivity such that $\lambda |\xi|^2 \leq a \xi \cdot \xi \leq |\xi|^2$

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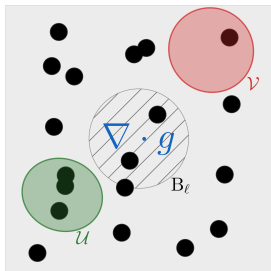
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Random heterogeneous media: a is sampled from a probability measure \mathbb{P}

- **Stationarity**, i.e. $a(\cdot + x)$ and a have the same distribution for all $x \in \mathbb{R}^3$

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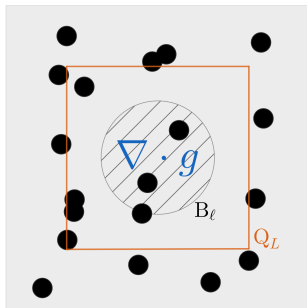
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Random heterogeneous media: a is sampled from a probability measure \mathbb{P}

- **Stationarity**, i.e $a(\cdot + x)$ and a have the same distribution for all $x \in \mathbb{R}^3$
- **Ergodicity**, i.e $a|_{\mathcal{U}}$ and $a|_{\mathcal{V}}$ decorrelates as $\text{dist}(\mathcal{U}, \mathcal{V}) \uparrow \infty$

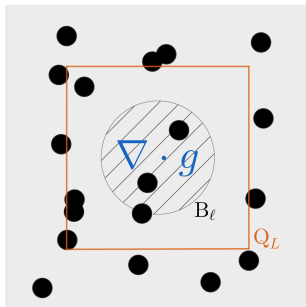
Computation of the electric field



In practice: We use a finite volume approximation

$$\begin{cases} -\nabla \cdot a \nabla u_L = \nabla \cdot g & \text{in } Q_L, \\ u_L = u_{bc} & \text{on } \partial Q_L. \end{cases}$$

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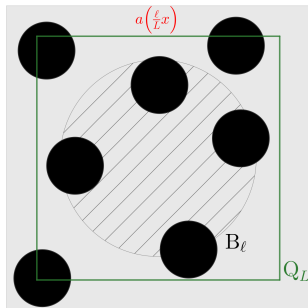
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Question: Construct u_{bc} that leads to optimal errors in L and ℓ !

Construction of the boundary condition

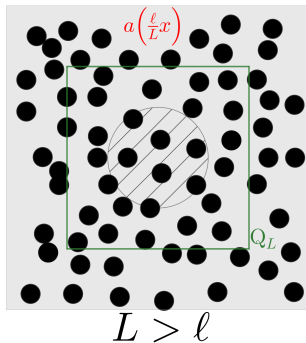
I. Large-scale behaviour in random media: Papanicolaou, Varadhan, Dal Maso, Modica; Gloria, Neukamm, Otto, Armstrong, Kuusi, Mourrat, Smart



$$L \sim \ell$$

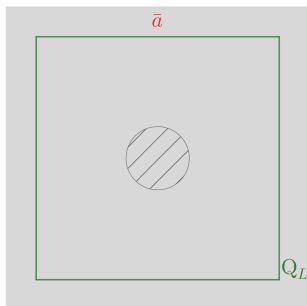
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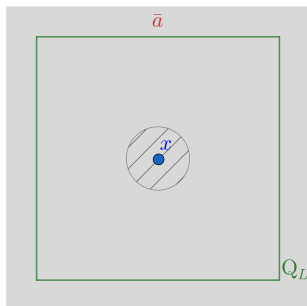
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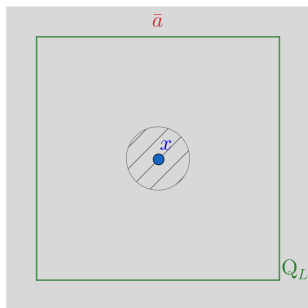
The media homogenizes and is described by its two-scale expansion

$$-\nabla \cdot a \nabla G(x, \cdot) = \delta(x - \cdot), \quad G(x, y) \approx \underbrace{(1 + \phi_i^{(1)}(x) \partial_{x_i})(1 + \phi_i^{(1)}(y) \partial_{y_i})}_{\text{Two-scale expansion}} \bar{G}(x - y)$$

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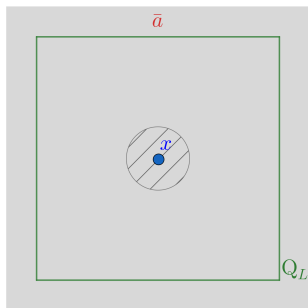
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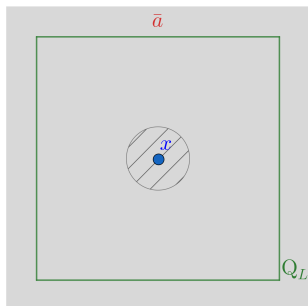
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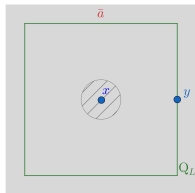
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II. Far-field behaviour in random media:

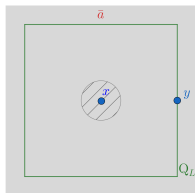


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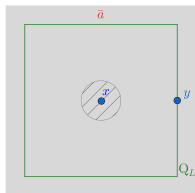
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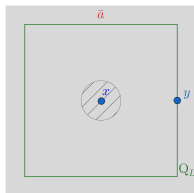
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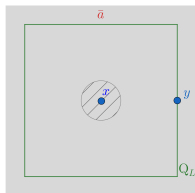
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$$u(x) \approx (1 + \phi_i^{(1)}(x) \partial_{x_i}) \left(\underbrace{\int_{\mathbb{R}^d} \bar{G}(x - y) \nabla \cdot g(y) dy}_{:= \bar{u}} - \int_{\mathbb{R}^d} \phi_i^{(1)}(y) \partial_i \bar{G}(x - y) \nabla \cdot g(y) dy \right).$$

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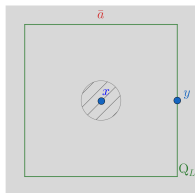
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In the regime $|x| = O(L) \gg |y| = O(\ell)$, we have $\partial_i \bar{G}(x - y) \approx \partial_i \bar{G}(x)$.

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Ansatz:

$$u_{bc} = (1 + \phi_i^{(1)} \partial_i) \left(\bar{u} - \partial_i \bar{G}(x) \int \phi_i^{(1)}(y) \nabla \cdot g(y) dy \right)$$

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Can we do better ? Depends how correlated is the medium.

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Theorem (Neukamm-Gloria-Otto ('20))

There exists $\phi^{(1)} \in H_{\text{loc}}^1$ and

$$\mathbb{E} \left[|\phi^{(1)}(x)|^2 \right]^{\frac{1}{2}} \lesssim \begin{cases} (1 + |x|)^{1 - \frac{\beta}{2}} & \text{for } \beta < 2 \\ \log^{\frac{1}{2}}(2 + |x|) & \text{for } \beta = 2 \\ 1 & \text{for } \beta > 2 \end{cases}$$

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For $\beta > 2$, we can go to a second-order approximation ! Fischer, Bella, Fehrman, Otto ('17)

A second-order expansion

Second-order correctors: Solution $\phi_{ij}^{(2)}$ of

$$-\nabla \cdot a \nabla \phi_{ij}^{(2)} = \nabla \cdot a \phi_i^{(1)} e_j + e_j \cdot (a(e_i + \nabla \phi_i^{(1)}) - \bar{a} e_i) \text{ in } \mathbb{R}^3$$

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Ansatz:

$$u_{bc} = \underbrace{(1 + \phi_i^{(1)} \partial_i + \phi_{ij}^{(2)} \partial_{ij})}_{2^{nd}\text{-order two-scale expansion}} \left(\bar{u} - \underbrace{\partial_i \bar{G}(x) \int \phi_i^{(1)}(y) \nabla \cdot g(y) dy}_{\text{Dipole correction}} + \underbrace{\partial_{ij} \bar{G}(x) c_{ij}}_{\text{Quadrupole correction}} \right)$$

The rigorous result

Theorem (C.-Wang ('24))

There exists $C, \gamma > 0$ such that for any $R \leq L$ and $L \geq C\ell$

$$\mathbb{P}\left(\left(\int_{B_R} |\nabla u - \nabla u_L|^2\right)^{\frac{1}{2}} \leq C \underbrace{\left(\frac{\ell}{L}\right)^d}_{\text{Dipole scaling}} \underbrace{L^{-\frac{\beta \wedge 3}{2}}}_{\text{Fluctuation scaling}}\right) \geq 1 - C \exp\left(-\ell^\gamma - R^\gamma - L^{\frac{1}{C}}\right)$$

Extend results from Jianfeng, Otto and Wang ('21) for the case $\beta \gg 1$

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How do we compute the boundary condition: There quantities to compute

$$\phi^{(1)}, \quad \phi^{(2)} \quad \text{and} \quad \bar{a}e_i := \mathbb{E}\left[a(e_i + \nabla\phi_i^{(1)})\right]$$

$$-\nabla \cdot a \nabla \phi_i^{(1)} = \nabla \cdot a e_i \text{ in } \mathbb{R}^3$$

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Approximation of the correctors

$$\bar{a}e_i = \mathbb{E} [a(e_i + \nabla\phi_i^{(1)})] = \lim_{L \uparrow \infty} \int_{Q_L} a(e_i + \nabla\phi_i^{(1)})$$

We solve

$$\frac{1}{T} \phi_{i,L}^{(1)} - \nabla \cdot a \nabla \phi_{i,L}^{(1)} = \nabla \cdot a e_i \text{ in } Q_L, \quad \phi_{i,L}^{(1)} = 0 \text{ on } \partial Q_L$$

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Theorem (C.-Wang ('24))

For any $p < \infty$

$$\mathbb{E} \left[\int_{Q_L} |(\phi_L^{(1)} - \phi^{(1)}, \sqrt{T}(\nabla \phi_L^{(1)} - \nabla \phi^{(1)}), \nabla \phi_L^{(2)} - \nabla \phi^{(2)})|^2 \right]^{\frac{1}{2}} \lesssim T^{\frac{1}{2} - \frac{\beta \wedge 3}{4}} + \left(\frac{\sqrt{T}}{L} \right)^p$$

References

The paper: C., Wang. *Artificial boundary conditions for random elliptic systems with correlated coefficient field*, Multiscale modelling and simulations ('24)

Related works:

- Jianfeng, Otto, Wang. *Optimal artificial boundary conditions based on second-order correctors for three dimensional random elliptic media*, Preprint ('21)
- C. *Optimal decay of the parabolic semigroup in stochastic homogenization for correlated coefficient fields*, Stochastics and Partial Differential Equations: Analysis and Computations ('22)
- Bella, Giunti, Otto. *Effective multipoles in random media*, Communications in Partial Differential Equations ('20)

Thank you for your attention