## Ensemble control of *n*-level systems via combined adiabatic and rotating wave approximations

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Let us consider a continuum of n-level systems described by the Schrödinger equation

$$i\dot{\psi}(t) = (H(\alpha) + \omega(t)H_c(\delta))\psi(t),$$
 (1)

where  $\omega(\cdot)$  is a real-valued control. Here we assume  $H(\alpha)$  is determined by an unknown parameter  $\alpha$  in a closed and connected domain  $\mathcal{D}$  of  $\mathbb{R}^m$ and has the following structure

$$H(\alpha) = \begin{pmatrix} \lambda_1(\alpha) & 0 & \dots & 0 \\ 0 & \lambda_2(\alpha) & 0 & \vdots \\ \vdots & 0 & \ddots & 0 \\ 0 & \dots & 0 & \lambda_n(\alpha) \end{pmatrix},$$

where  $\lambda_j : \mathcal{D} \to \mathbb{R}$  is a continuous function for each  $j \in \{1, \dots, n\}$ .

 $H_c(\delta)$  is a self-adjoint matrix that describes the control coupling between the eigenstates of the system and has the form

$$H_c(\delta) = \begin{pmatrix} 0 & \delta_{12} & \dots & \delta_{1n} \\ \delta_{12} & 0 & \ddots & \vdots \\ \vdots & \ddots & \ddots & \delta_{n-1,n} \\ \delta_{1n} & \cdots & \delta_{n-1,n} & 0 \end{pmatrix},$$

We assume that each  $\delta_{jk}$  is unknown but that it belongs to some closed interval  $\mathcal{I}_{jk} = \left[\delta_{jk}^{0}, \delta_{jk}^{1}\right]$  in  $\mathbb{R}$  such that  $0 \notin \mathcal{I}_{jk}$ .

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For  $p, q \in \{1, ..., n\}$  s.t.  $p \neq q$ , one would like to find a uniform control  $\omega(\cdot)$  for the family of quantum systems s.t. if at t = 0, all systems are at the same initial state  $\psi(0) = \mathbf{e}_p$ , then at time T the systems are close to final states of form  $e^{i\theta}\mathbf{e}_q$  for some  $\theta \in \mathbb{R}$ .

#### Two-level system with complex control

Consider a two-level system with a complex control

$$i\frac{\mathsf{d}}{\mathsf{d}t}\psi(t) = \begin{pmatrix} E+\alpha & \omega(t)\\ \omega^*(t) & -E-\alpha \end{pmatrix}\psi(t),\tag{2}$$

where  $\alpha$  is in an unknown parameter in an interval  $[\alpha_0, \alpha_1]$ . For  $\epsilon > 0$ , choose the control

$$\omega_{\epsilon}(t) = u(\epsilon t) e^{-i\left(2Et + \frac{\Delta(\epsilon t)}{\epsilon}\right)},$$
(3)

where u and  $\Delta$  are real-valued functions on [0, T]. Let us apply the change of variables

$$\psi(t) = \begin{pmatrix} e^{-i\left(Et + \frac{\Delta(\epsilon t)}{2\epsilon}\right)} & 0\\ 0 & e^{i\left(Et + \frac{\Delta(\epsilon t)}{2\epsilon}\right)} \end{pmatrix} \Psi(t).$$
(4)

Then we can obtain that

$$i\frac{\mathsf{d}}{\mathsf{d}t}\Psi = \begin{pmatrix} \alpha - v(\epsilon t) & u(\epsilon t) \\ u(\epsilon t) & -\alpha + v(\epsilon t) \end{pmatrix}\Psi,\tag{5}$$

where  $v(s) = \Delta'(s)/2$ . Assume that

$$\sqrt{(\alpha - \nu(s))^2 + u(s)^2} > 0, \quad \forall s \in [0, T].$$
(6)

Then by adiabatic theorem (e.g. [Teufel, 2003]), when  $\epsilon \to 0$ , if  $\Psi(0)$  is in the eigenspace associated with  $\lambda_+(0)$  (respectively,  $\lambda_-(0)$ ), then the state will remain close to the eigenspace associated with  $\lambda_+(t)$  (respectively,  $\lambda_-(t)$ ).

#### Two-level system with complex control

Fix T = 1. Define  $u(\cdot)$  and  $v(\cdot)$  as follow



Therefore, the solution  $\psi$  of (2) with initial state (0,1) and corresponding to  $\omega_{\epsilon}$  satisfies  $|\psi(T/\epsilon) - (e^{i\theta}, 0)| \leq C\epsilon$  for some  $\theta \in \mathbb{R}$ .

## Two-level system with real control

Consider a two-level system with a real-valued control

$$i\frac{d}{dt}\psi(t) = \begin{pmatrix} E+\alpha & \omega(t)\\ \omega(t) & -E-\alpha \end{pmatrix}\psi(t),$$
(7)

and the real chirped pulse

$$\omega_{\epsilon}(t) = 2u(\epsilon t)\cos\left(2Et + \Delta(\epsilon t)/\epsilon\right),\tag{8}$$

and apply the change of variables from (4). We obtain

$$i\frac{d}{dt}\Psi = \begin{bmatrix} \begin{pmatrix} \alpha - v(\epsilon t) & u(\epsilon t) \\ u(\epsilon t) & -\alpha + v(\epsilon t) \end{pmatrix} + R(t,\epsilon) \end{bmatrix} \Psi,$$
(9)  
$$R(t,\epsilon) = \begin{pmatrix} 0 & u(\epsilon t)e^{i(4Et+2\Delta(\epsilon t)/\epsilon)} \\ u(\epsilon t)e^{-i(4Et+2\Delta(\epsilon t)/\epsilon)} & 0 \end{pmatrix}.$$

The adiabatic theorem discussed above is no longer valid.

# RWA(Rotating wave approximations)

If  $\alpha = 0$ , the dynamics of a system with complex-valued control can be simulated with a real-valued control. For  $\epsilon > 0$ , consider the control

$$\omega_{\epsilon}(t) = 2\epsilon u(\epsilon t) \cos(2Et + \Delta(\epsilon t)). \tag{10}$$

Applying the change of variable

$$\psi(t) = \begin{pmatrix} e^{-i(Et + \Delta(\epsilon t)/2)} & 0\\ 0 & e^{i(Et + \Delta(\epsilon t)/2)} \end{pmatrix} \hat{\psi}(t), \quad (11)$$

and the change of variable  $ilde{\psi}(s) = \hat{\psi}(\epsilon t)$ , we obtain

$$i\frac{\mathsf{d}}{\mathsf{d}s}\tilde{\psi} = \begin{bmatrix} \begin{pmatrix} -v(s) & u(s) \\ u(s) & v(s) \end{pmatrix} + \begin{pmatrix} 0 & e^{i(4Es/\epsilon + 2\Delta(s))}u(s) \\ e^{-i(4Es/\epsilon + 2\Delta(s))}u(s) & 0 \end{pmatrix} \end{bmatrix} \tilde{\psi}$$

By classical averaging result, the second term induces a perturbation of order  $\epsilon$  to the solution.

## RWA(Rotating wave approximations)

However, when the resonance is not precisely known  $\alpha \in [\alpha_0, \alpha_1]$ ,

$$irac{\mathsf{d}}{\mathsf{d}s} ilde{\psi}_lpha = egin{bmatrix} lpha/\epsilon - oldsymbol{v}(s) & u(s) \ u(s) & -lpha/\epsilon + v(s) \end{pmatrix} + B(s,\epsilon) \end{bmatrix} ilde{\psi}_lpha,$$

where  $B(s, \epsilon)$  denotes the second term in the previous equation. The adiabatic following is no longer valid since  $\lim_{\epsilon \to 0} \alpha_1/\epsilon = +\infty$  and  $\lim_{\epsilon \to 0} \alpha_0/\epsilon = -\infty$ .

Combination of RWA and AA ( Robin, Augier, Boscain, Sigalotti, 2022)

Consider two time scales  $\epsilon_1, \epsilon_2 > 0$  and a control law of the type

$$\omega_{\epsilon_1,\epsilon_2}(t) = 2\epsilon_1 \delta u(\epsilon_1 \epsilon_2 t) \cos\left(2Et + \frac{\Delta(\epsilon_1 \epsilon_2 t)}{\epsilon_1 \epsilon_2}\right).$$
(12)

Here  $u, \Delta$  are real-valued smooth functions on [0,T] and  $\delta \in \mathbb{R}^*$  is an unknown parameter.

#### Theorem (Robin, Augier, Boscain, Sigalotti, 2022)

Assume that  $v_0 < 0 < v_1$  are such that  $3(E + v_0) \ge E + v_1$ . Fix T > 0and  $u, \Delta : [0, T] \to \mathbb{R}$  smooth (e.g.,  $u \in C^2$  and  $\Delta \in C^3$ ) such that

- **(** $u(0), \Delta'(0)$  $) = (0, 2v_0)$  and  $(u(T), \Delta'(T)) = (0, 2v_1);$
- 2  $\forall s \in (0, T), u(s) > 0 \text{ and } \Delta''(s) \ge 0.$

Denote by  $\psi_{\epsilon_1,\epsilon_2}^{\alpha}$  the solution of (7) with initial condition  $\psi_{\epsilon_1,\epsilon_2}^{\alpha}(0) = (0,1)$ and control  $\omega_{\epsilon_1,\epsilon_2}$  as in (12). Then for every  $N_0 \in \mathbb{N}$ , for every compact interval  $I \subseteq (v_0, v_1)$ , there exist  $C_{N_0} > 0$  and  $\eta > 0$  such that for every  $\alpha \in I$  and every  $(\epsilon_1, \epsilon_2) \in (0, \eta)^2$ ,

$$\left|\psi_{\epsilon_1,\epsilon_2}(T/(\epsilon_1\epsilon_2)) - (e^{i\theta},0)\right| < C_{N_0}\max(\epsilon_2/\epsilon_1,\epsilon_1^{N_0-1}/\epsilon_2)$$
 (13)

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for some  $\theta \in \mathbb{R}$ . Moreover, the constant  $C_{N_0}$  can be taken locally uniform with respect to the parameter  $\delta > 0$  appearing in (12).

### Combination of RWA and AA in population inversion

The conditions in the theorem ensures that we can simulate the dynamics

$$\frac{\mathsf{d}}{\mathsf{d}t}\psi_{\mathsf{slow}} = \begin{pmatrix} \mathsf{v}(\epsilon_1\epsilon_2t) - \alpha & \epsilon_1 u(\epsilon_1\epsilon_2t) \\ \epsilon_1 u(\epsilon_1\epsilon_2t) & -\mathsf{v}(\epsilon_1\epsilon_2t) + \alpha \end{pmatrix} \psi_{\mathsf{slow}}.$$
 (14)

A non-classical adiabatic theorem can be applied with an error  $\mathcal{O}(\epsilon_2/\epsilon_1)$ .

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#### Recall the *n*-level system

$$i\frac{\mathsf{d}}{\mathsf{d}t}\psi(t) = \begin{pmatrix} \lambda_1(\alpha) & \delta_{1,2}\omega(t) & \dots & \delta_{1,n}\omega(t) \\ \delta_{1,2}\omega(t) & \lambda_2(\alpha) & \ddots & \vdots \\ \vdots & \ddots & \ddots & \delta_{n-1,n}\omega(t) \\ \delta_{1,n}\omega(t) & \dots & \delta_{n-1,n}\omega(t) & \lambda_n(\alpha) \end{pmatrix} \psi(t),$$

and the control law

$$\omega_{\epsilon_1\epsilon_2}(t) = 2\epsilon_1 u(\epsilon_1\epsilon_2 t) \cos\left(\int_0^t f(\epsilon_1\epsilon_2 \tau) \mathrm{d}\tau\right),\tag{15}$$

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where  $u, f : [0, T] \rightarrow \mathbb{R}$  are functions to be chosen.

Let us assume that for all  $1 \le j < k \le n$ , and for all  $\alpha \in D$ ,  $\lambda_k(\alpha) - \lambda_j(\alpha) > 0$ . Fix  $1 \le p < q \le n$  and assume that there exist  $0 < v_0 < v_1$  such that

• For  $\alpha \in \mathcal{D}$ ,  $\lambda_p(\alpha) - \lambda_q(\alpha) \in (v_0, v_1)$  and, for all  $1 \le j < k \le n$  such that  $(j, k) \ne (p, q)$ , we have that  $\forall \alpha \in \mathcal{D}$ ,  $\lambda_k(\alpha) - \lambda_j(\alpha) \notin [v_0, v_1]$ .

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**2** For all  $1 \le j < k \le n$  and all  $\alpha \in \mathcal{D}$ ,  $\lambda_k(\alpha) - \lambda_j(\alpha) \notin [2v_0, 2v_1]$ .

Then we can fix T > 0 and take  $u, f \in C^2([0, T], \mathbb{R})$  such that

i) 
$$(u(0), f(0)) = (0, v_0)$$
 and  $(u(T), f(T)) = (0, v_1);$ 

ii) 
$$\forall s \in (0, T), u(s) > 0, f'(s) > 0.$$

Denote by  $\psi_{\epsilon_1,\epsilon_2}$  the solution of (1) with initial condition  $\psi_{\epsilon_i,\epsilon_2}(0) = \mathbf{e}_p$ and the control law  $\omega_{\epsilon_1,\epsilon_2}$  as in (15). Then there exist C > 0 and  $\eta > 0$ such that for every  $\alpha \in \mathcal{D}$  and every  $(\epsilon_1,\epsilon_2) \in (0,\eta)^2$ ,

$$\left\|\psi_{\epsilon_{1},\epsilon_{2}}\left(\frac{T}{\epsilon_{1}\epsilon_{2}}\right) - \exp(i\theta)\mathbf{e}_{q}\right\| < C \max\left(\frac{\epsilon_{1}^{2}}{\epsilon_{2}},\frac{\epsilon_{2}}{\epsilon_{1}}\right)$$
(16)

for some  $\theta \in \mathbb{R}$ .

For  $E \in \mathbb{R}$  and  $1 \leq j \leq k \leq n$ , let us define

$$A_{jk}(E) = \begin{cases} \exp(iE)\mathbf{e}_{jk} + \exp(-iE)\mathbf{e}_{kj} & \text{if } j < k, \\ \cos(E)\mathbf{e}_{jj} & \text{if } j = k, \end{cases}$$
$$B_{jk}(E) = \begin{cases} i \exp(iE)\mathbf{e}_{jk} - i \exp(-iE)\mathbf{e}_{kj} & \text{if } j < k, \\ -\sin(E)\mathbf{e}_{jj} & \text{if } j = k. \end{cases}$$

For  $1 \leq j < k \leq n$  and  $\sigma \in \mathbb{Z}$ ,  $\phi^{\sigma}_{jk}(t)$  is defined by

$$\phi_{jk}^{\sigma}(t) = (\lambda_j - \lambda_k)t + \sigma \int_0^t f(\epsilon_1 \epsilon_2 \tau) \mathrm{d}\tau.$$

Let us recast (1) in the interaction frame  $\psi(t) = \exp(-itH(\alpha))\psi_I(t)$ . The dynamics of  $\psi_I$  is characterized by the Hamiltonian

$$H_{I}(t) = \sum_{j=1}^{n-1} \sum_{k=j+1}^{n} \delta_{j,k} u(\epsilon_{1}\epsilon_{2}t) \left[ A_{jk}(\phi_{jk}^{1}(t)) + A_{jk}(\phi_{jk}^{-1}(t)) \right].$$

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Assume that hypothesis 1 of Theorem is satisfied. Then there exist a function  $X_1(t)$  on  $[0, 1/(\epsilon_1 \epsilon_2)]$  and a change of variables

$$\psi_I(t) = \exp(i\epsilon_1 X_1(t))\hat{\psi}_1(t), \qquad (17)$$

such that the dynamics of  $\hat{\psi}_1$  is characterized by the Hamiltonian

$$\begin{split} \hat{H}_{1}(t) &= \epsilon_{1} \delta_{pq} u(\epsilon_{1} \epsilon_{2} t) A_{pq} \left( \phi_{pq}^{1}(t) \right) + \sum_{(j,k,\sigma) \in \mathcal{J}} \epsilon_{1}^{2} h_{jk}^{\sigma}(\epsilon_{1} \epsilon_{2} t) A_{jk} \left( \phi_{jk}^{\sigma}(t) \right) \\ &+ \mathcal{O}(\epsilon_{1}^{3} + \epsilon_{1}^{2} \epsilon_{2}), \end{split}$$

where  $\mathcal{J} = \{(j, k, \sigma) \mid 1 \leq j \leq k \leq n, \sigma \in \{-2, 0, 2\}\}$  and  $h_{jk}^{\sigma}(\cdot)$  are functions on [0, 1].

Similar to the reasoning above, with Hypothesis 2, there exist a function  $X_2(t)$  on  $[0, 1/(\epsilon_1 \epsilon_2)]$  and a change of variables

$$\hat{\psi}_1(t) = \exp(i\epsilon_1^2 X_2(t))\hat{\psi}_2(t),$$

such that the dynamics of  $\hat{\psi}_2(t)$  is characterized by the Hamiltonian

$$\begin{split} \hat{H}_2(t) &= \hat{H}_1(t) + \epsilon_1^2 \frac{\mathsf{d}}{\mathsf{d}t} X_2(t) + \mathcal{O}(\epsilon_1^3) \\ &= \epsilon_1 \delta_{pq} u(\epsilon_1 \epsilon_2 t) A_{pq} \left( \phi_{pq}^1(t) \right) + \sum_{j=1}^n \epsilon_1^2 h_{jj}^0(\epsilon_1 \epsilon_2 t) \mathbf{e}_{jj} + \mathcal{O}(\epsilon_1^3 + \epsilon_1^2 \epsilon_2). \end{split}$$

Since u(0) = u(1) = 0, we deduce that  $h_{jk}^{\sigma}(0) = h_{jk}^{\sigma}(1) = 0$  for every  $(j, k, \sigma) \in \mathcal{J}$ . Then  $X_1(0) = X_1\left(\frac{1}{\epsilon_1\epsilon_2}\right) = 0$  and  $X_2(0) = X_2\left(\frac{1}{\epsilon_1\epsilon_2}\right) = 0$ . Hence,  $\psi_I(0) = \hat{\psi}_1(0) = \hat{\psi}_2(0)$  and  $\psi_I\left(\frac{1}{\epsilon_1\epsilon_2}\right) = \hat{\psi}_1\left(\frac{1}{\epsilon_2\epsilon_2}\right) = \hat{\psi}_2\left(\frac{1}{\epsilon_2\epsilon_2}\right)$ .

Assume that  $\psi_I(0) \in V = \operatorname{span}(\mathbf{e}_p, \mathbf{e}_q)$ , let us first project  $\hat{H}_2(t)$  on V and introduce the following truncation

$$H_{rwa}(t) = \epsilon_1 \delta_{pq} u(\epsilon_1 \epsilon_2 t) A_{pq} (\phi_{pq}^1(t)) + \sum_{j \in \{p,q\}} \epsilon_1^2 h_{jj}^0 (\epsilon_1 \epsilon_2 t) \mathbf{e}_{jj},$$

and the system

$$irac{\mathsf{d}}{\mathsf{d}t}\psi_\mathsf{rwa}(t)=H_\mathsf{rwa}(t)\psi_\mathsf{rwa}(t), \quad \psi_\mathsf{rwa}(0)=\psi_I(0).$$

We can prove that

$$\|\psi_I\left(1/(\epsilon_1\epsilon_2)\right) - \psi_{\mathsf{rwa}}\left(1/(\epsilon_1\epsilon_2)\right)\| = \mathcal{O}(\epsilon_1^2/\epsilon_2 + \epsilon_1).$$

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Introduce the unitary change of variables with  $\psi_{\sf rwa}(t) = U(t)\psi_{\sf slow}(t)$  with

$$U(t) = e_{\rho\rho} + \exp\left(-i\phi_{\rho q}^{1}(t)\right) \mathbf{e}_{qq}.$$
(18)

The dynamics of  $\psi_{\rm slow}$  is characterized by the Hamiltonian

$$egin{aligned} \mathcal{H}_{\mathsf{slow}}(t) &= \epsilon_1^2 h_{
hop}^0(\epsilon_1 \epsilon_2 t) \mathbf{e}_{
hop} + \epsilon_1 \delta_{
hop} u(\epsilon_1 \epsilon_2 t) (\mathbf{e}_{
hop} + \mathbf{e}_{q
ho}) \ &+ ig( (\lambda_2 - \lambda_1 - f(\epsilon_1 \epsilon_2 t) + \epsilon_1^2 h_{qq}^0(\epsilon_1 \epsilon_2 t) ig) \mathbf{e}_{qq}. \end{aligned}$$

#### Remark

It has been shown in [Robin, Augier, Boscain, Sigalotti, 2022] that the diagonal terms  $\epsilon_1^2 h_{qq}^0(\epsilon_1 \epsilon_2 t) \mathbf{e}_{qq}$  and  $\epsilon_1^2 h_{pp}^0(\epsilon_1 \epsilon_2 t) \mathbf{e}_{pp}$  can be neglected when applying the adiabatic theorem.



We can then apply Lemma 30 in [Robin, Augier, Boscain, Sigalotti, 2022] and obtain that, for the control  $\omega_{\epsilon_1,\epsilon_2}(\cdot)$  given in (15), there exists C' > 0 such that for  $\psi_{\mathsf{rwa}}(0) = \mathbf{e}_p$ 

$$\min_{\theta \in [0,2\pi]} \left\| \psi_{\mathsf{rwa}} \left( 1/(\epsilon_1 \epsilon_2) \right) - \exp(i\theta) \mathbf{e}_q \right\| \le C' \epsilon_2 / \epsilon_1.$$

We deduce that, in the interaction frame, there exists C > 0 such that

$$\min_{\theta \in [0,2\pi]} \left\| \psi_I\left(\frac{1}{\epsilon_1 \epsilon_2}\right) - \exp(i\theta) \mathbf{e}_q \right\| \leq C \max\left(\frac{\epsilon_1^2}{\epsilon_2}, \frac{\epsilon_2}{\epsilon_1}\right).$$

It follows that

$$\min_{\theta \in [0,2\pi]} \left\| \psi_{\epsilon_1,\epsilon_2} \left( \frac{1}{\epsilon_1 \epsilon_2} \right) - \exp(i\theta) \mathbf{e}_q \right\| < C \max\left( \frac{\epsilon_1^2}{\epsilon_2}, \frac{\epsilon_2}{\epsilon_1} \right).$$

We can thus conclude the proof.

Consider a system as in (1) with drift and control Hamiltonians

$$H(\alpha) = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1.5 + \alpha & 0 \\ 0 & 0 & 8 \end{pmatrix}, \quad H_c = \begin{pmatrix} 0 & 1 & 5 \\ 1 & 0 & 5 \\ 5 & 5 & 0 \end{pmatrix}.$$

In order to realize a population inversion between the first and the second eigenstates, we fix T = 1,  $(\epsilon_1, \epsilon_2) = (10^{-2}, 10^{-7/2})$ ,  $\psi_{\epsilon_1, \epsilon_2}(0) = (1, 0, 0)^{\top}$  and use the control law given in (15) with  $v_0 = 1$ ,  $v_1 = 2.5$ . The fidelity at  $s \in [0, 1]$  is defined as

$$\operatorname{fid}(s) = 1 - \min_{\theta \in [0, 2\pi]} \left\| \psi_{\epsilon_1, \epsilon_2} \left( \frac{s}{\epsilon_1 \epsilon_2} \right) - e^{i\theta} \mathbf{e}_q \right\|.$$
(19)

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Here, the target state  $\mathbf{e}_q = (0, 1, 0)^{\top}$ .

Example



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