

**NUMERICAL METHODS  
FOR THE  
FREE-SURFACE NAVIER-STOKES  
EQUATIONS  
AND APPLICATIONS TO  
BREAKING WATER WAVES**

**CANUM 2024**  
MODÉLISATION, MÉTHODES NUMÉRIQUES ET APPLICATIONS EN  
OCÉANOGRAPHIE

**Alan Riquier** - Département de Mathématiques et Applications (ENS - PSL)

Joint work with Emmanuel Dormy (DMA - ENS PSL)

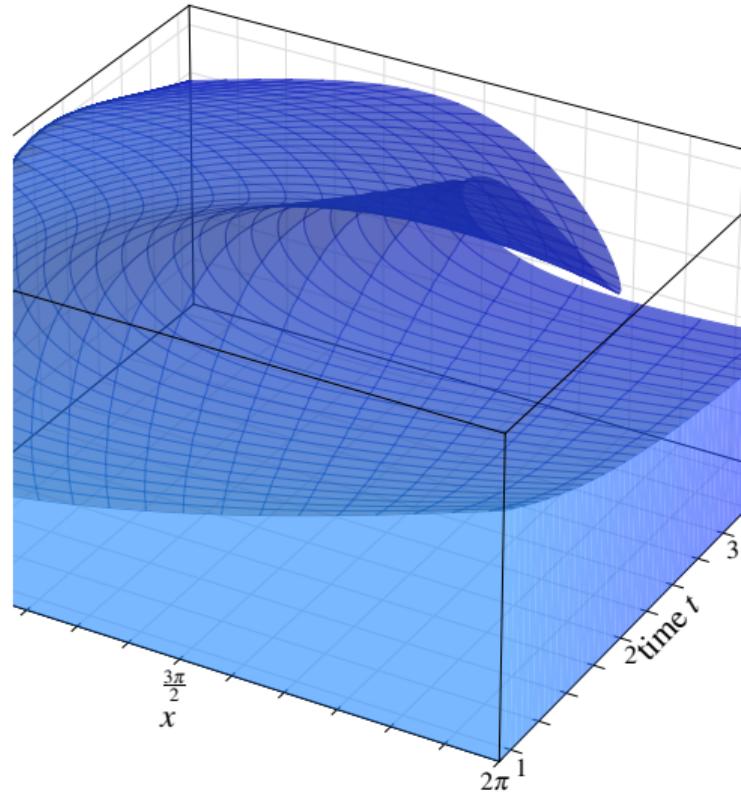
May 28, 2024



# Outline

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1. The free-surface Navier-Stokes equations
2. Finite element framework
3. The  $\text{Re} \rightarrow +\infty$  limit
4. The superficial boundary layer



# Problem formulation

Navier-Stokes

Lagrangian advection

Initial condition



# Viscous Water Waves

## Nondimensionalization

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Nondimensional quantities are defined as follows

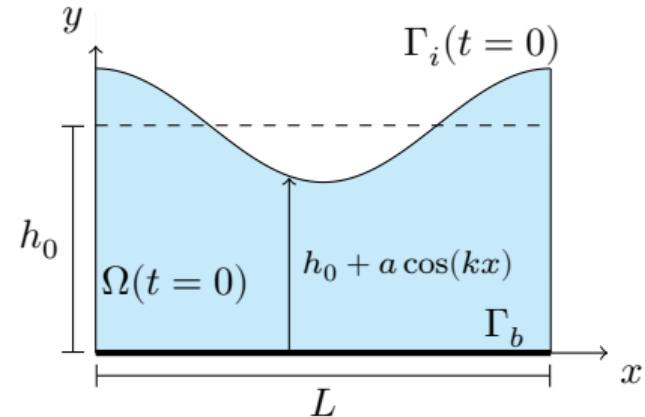
$$x \rightarrow h_0 x$$

$$u \rightarrow \sqrt{gh_0} \cdot u$$

$$p \rightarrow \rho gh_0 \cdot p$$

This allows to define the **Reynolds number**  $\text{Re}$ ,

$$\text{Re} = \frac{\rho h_0 \sqrt{gh_0}}{\mu}$$



# Viscous Water Waves

Navier-Stokes equation

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Incompressible, non-dimensional, **Navier-Stokes** equation in  $\Omega(t)$ :

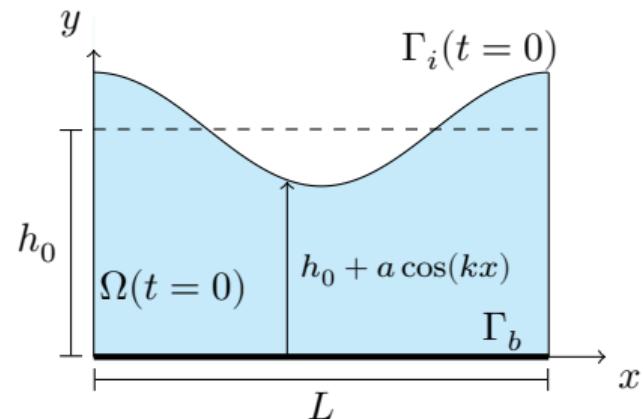
$$\begin{cases} \partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p + \frac{1}{\text{Re}} \Delta \mathbf{u} + \mathbf{g} \\ \nabla \cdot \mathbf{u} = 0 \end{cases}$$

Navier boundary conditions on  $\Gamma_b$ ,

$$\mathbf{u} \cdot \mathbf{n} = 0 \quad ; \quad \mathbf{t} \cdot [\nabla \mathbf{u} + (\nabla \mathbf{u})^t] \cdot \mathbf{n} = 0$$

Stress-free boundary condition on  $\Gamma_i(t)$ ,

$$p \mathbf{n} - \frac{1}{\text{Re}} \cdot [\nabla \mathbf{u} + (\nabla \mathbf{u})^t] \cdot \mathbf{n} = 0$$



# Weak formulation

Navier-Stokes problem

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## Function space

$$\mathbf{H}_{\Gamma_b}^1(\Omega(t)) = \left\{ \mathbf{v} \in H^1(\Omega(t); \mathbb{R}^2) : \mathbf{v} \cdot \mathbf{n} = 0 \text{ on } \Gamma_b \right\}$$

We **do not** assume **the** incompressibility,  $\nabla \cdot \mathbf{u} = 0$ , directly in the function space as it would not work in finite element!

Find  $\mathbf{u} \in \mathcal{C}^1([0, T); \mathbf{H}_{\Gamma_b}^1)$  and  $p \in L^\infty([0, T), L^2)$  such that

$$\int_{\Omega(t)} \mathbf{v} \cdot \partial_t \mathbf{u} + \mathbf{v} \cdot (\mathbf{u} \cdot \nabla) \mathbf{u} + \frac{2}{\text{Re}} \mathbb{S}(\mathbf{v}) : \mathbb{S}(\mathbf{u}) - p \nabla \cdot \mathbf{v} + q \nabla \cdot \mathbf{u} - \mathbf{v} \cdot \mathbf{g} = 0$$

for all  $\mathbf{v} \in \mathbf{H}_{\Gamma_b}^1$  and  $q \in L^2$ , at all time  $t \in (0, T)$ .

# Surface tension

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Let Bo the **Bond number**,

$$\text{Bo} = \frac{\rho g L^2}{\sigma}$$

and  $\kappa : [0, T) \times \Gamma_s(t) \rightarrow \mathbb{R}$  the **surface curvature**.

Find  $\mathbf{u} \in \mathcal{C}^1([0, T); \mathbf{H}_{\Gamma_b}^1)$  and  $p \in L^\infty([0, T), L^2)$  such that

$$\int_{\Omega(t)} \mathbf{v} \cdot \partial_t \mathbf{u} + \mathbf{v} \cdot (\mathbf{u} \cdot \nabla) \mathbf{u} + \frac{2}{\text{Re}} \mathbb{S}(\mathbf{v}) : \mathbb{S}(\mathbf{u}) - p \nabla \cdot \mathbf{v} + q \nabla \cdot \mathbf{u} - \mathbf{v} \cdot \mathbf{g} = \int_{\Gamma_s(t)} \kappa \text{Bo}^{-1} \mathbf{v} \cdot \mathbf{n} \, dS$$

for all  $\mathbf{v} \in \mathbf{H}_{\Gamma_b}^1$  and  $q \in L^2$ , at all time  $t \in (0, T)$ .

# Interface advection

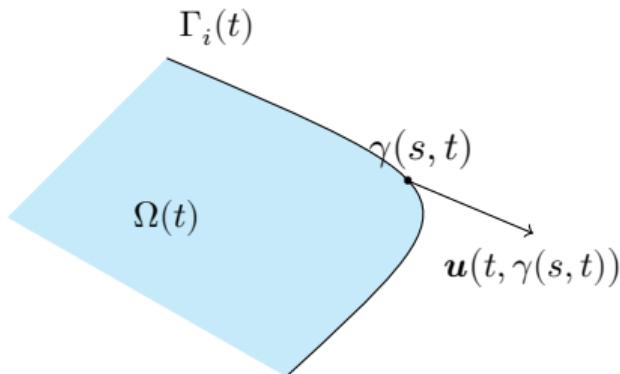
Lagrangian scheme

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Interface is a **parametrised curve**  $\gamma(s, t) \in \mathbb{R}^2$  whose evolution is given by

$$\frac{\partial \gamma}{\partial t}(s, t) = \mathbf{u}(t, \gamma(s, t))$$

i.e. points on the interface have the **same velocity** as the fluid particles.



The interface contains all points parametrised by  $\gamma$ ,

$$\Gamma_i(t) = \bigcup_s \{\gamma(s, t)\}$$

# Initial conditions

Theory of linear waves

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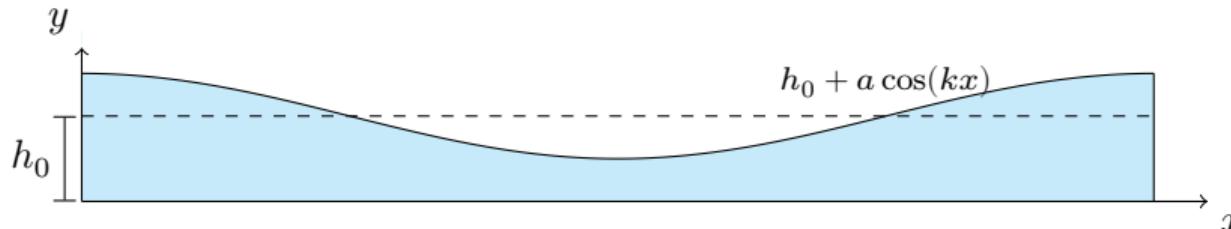
Linear wave of (small) amplitude  $a$ ,

$$\gamma_0(t, x) = h_0 + a \cos(kx - \omega t) \quad \text{with} \quad \omega = \sqrt{gk \tanh(kh_0)}$$

$$\phi_0(t, x, y) = \frac{a\omega}{k} \frac{\cosh(ky)}{\sinh(kh_0)} \cdot \sin(kx - \omega t) + \mathcal{O}(ka)$$

The velocity is then

$$\mathbf{u}_0(t, x, y) = \nabla \phi_0 = \frac{a\omega}{\sinh(kh_0)} \cdot \begin{bmatrix} \cosh(ky) \cos(kx - \omega t) \\ \sinh(ky) \sin(kx - \omega t) \end{bmatrix}$$



# Initial conditions

Laplace problem

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$$\mathbf{u}_0(t, x, y) = \frac{a\omega}{\sinh(kh_0)} \cdot \begin{bmatrix} \cosh(ky) \cos(kx - \omega t) \\ \sinh(ky) \sin(kx - \omega t) \end{bmatrix}$$

so the velocity  $\mathbf{n} \cdot \mathbf{u}_0$  in the **normal** direction can be computed from

$$\mathbf{n}_0(x) = \frac{1}{\sqrt{1 + (\partial_x \gamma_0)^2}} \begin{bmatrix} -\partial_x \gamma_0 \\ 1 \end{bmatrix} \quad \text{where } \gamma_0(x) = h_0 + a \cos(kx)$$

The **initial velocity** is then constructed from assuming  $y = h_0$  in  $\mathbf{u}_0$  and solving the following Laplace problem

$$\begin{cases} \Delta\phi = 0 & \text{in the domain } \Omega(t) \\ \partial_n \phi = 0 & \text{at the bottom } \Gamma_b \\ \partial_n \phi = \mathbf{n}_0 \cdot \mathbf{u}_0(y = h_0) & \text{at the surface } \Gamma_i(t = 0) \end{cases}$$

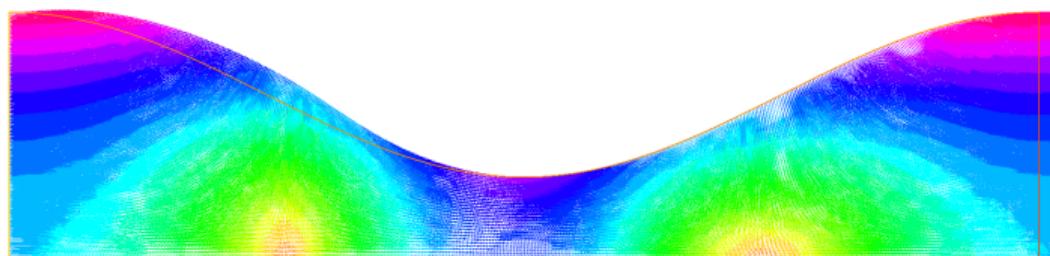
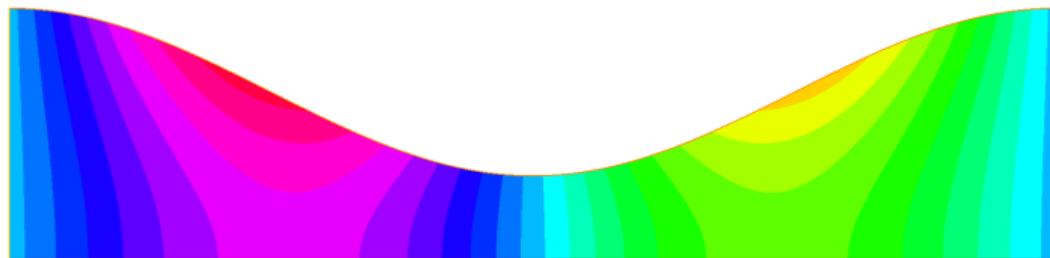
and the **initial velocity** is  $\mathbf{u}(t = 0) = \nabla\phi$ .

# Initial condition

Prescribed normal velocity

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$$\partial_n \phi = \mathbf{n}_0 \cdot \mathbf{u}_0(y = h_0)$$



Large values

Small values

# Numerical framework

Finite elements

Mixed Lagrangian-Eulerian

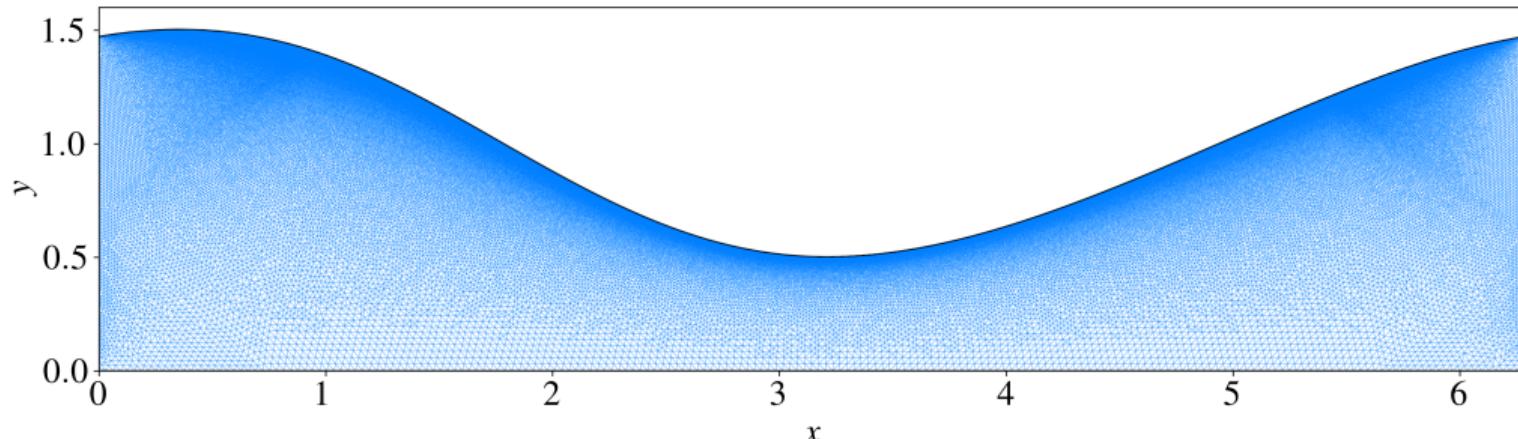
Geometric multigrid solver



# Finite Elements discretization

We use the FreeFem finite elements library  [Hecht \(2012\)](#) for

- Mesh generation and handling
- Matrices computations and handling
- Interface with PETSc



4 000 points on the interface, initially  $\approx 200\,000$  triangles,  $\approx 10^6$  degrees of freedom.

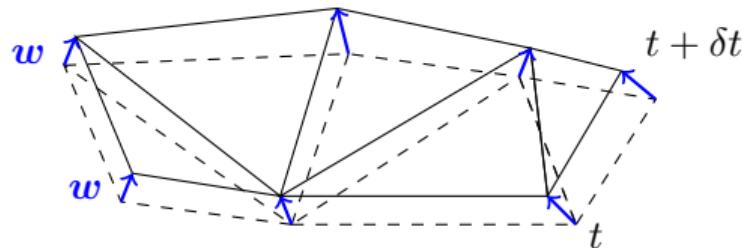
# Mesh advection scheme

Let  $w$  the **velocity of the mesh**. At each time step, we numerically solve the problem

$$\begin{cases} \Delta w = 0 & \text{in } \Omega_t \\ w = u & \text{on } \Gamma_{s,t} \\ w = 0 & \text{on } \Gamma_b \end{cases}$$

And each point of the mesh is advected with velocity  $w$ . Points on the interface are thus **purely Lagrangian!**

This is called the **Arbitrary Lagrangian Eulerian** method (ALE).



# Time stepping scheme

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Crank-Nicolson second order in time scheme. CFL condition to compute the time step at each iteration.

At each time step, we solve the following problem

$$\begin{aligned} \int_{\Omega(t)} \mathbf{v} \cdot \frac{\mathbf{u}^{n+1} - \mathbf{u}^n}{\delta t} + \mathbf{v} \cdot ((\mathbf{u}^n - \mathbf{w}^n) \cdot \nabla) \left( \frac{\mathbf{u}^{n+1} + \mathbf{u}^n}{2} \right) + \frac{2}{\text{Re}} \mathbb{S}(\mathbf{v}) : \mathbb{S} \left( \frac{\mathbf{u}^{n+1} + \mathbf{u}^n}{2} \right) \\ + \left( \frac{p^{n+1} + p^n}{2} \right) \nabla \cdot \mathbf{v} + q \nabla \cdot \left( \frac{\mathbf{u}^{n+1} + \mathbf{u}^n}{2} \right) - \mathbf{v} \cdot \mathbf{g} \\ = 0 \end{aligned}$$

for all  $(\mathbf{v}, q)$  and then compute  $\mathbf{w}^{n+1}$  before advecting the mesh.

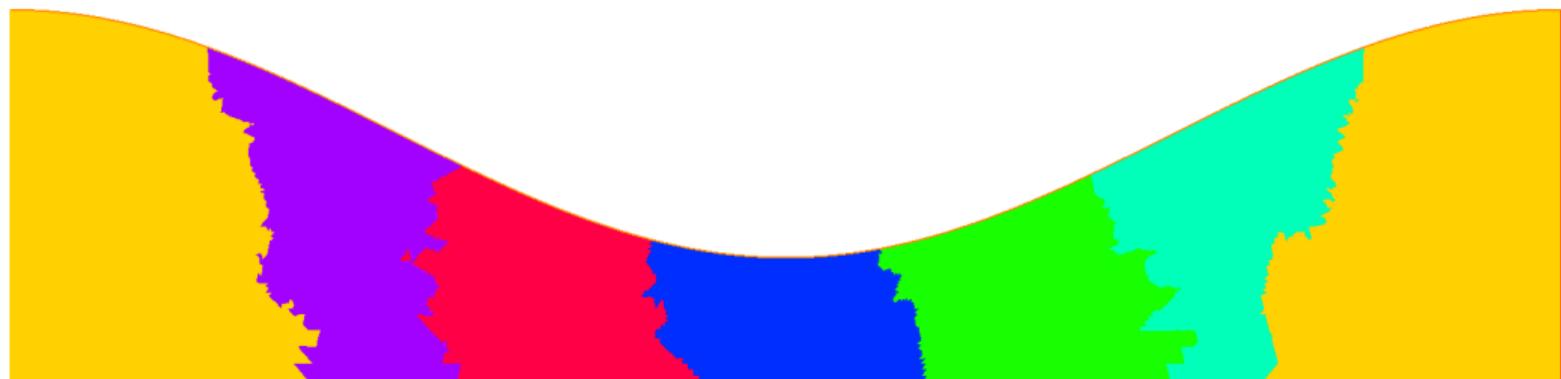
We use  $\mathbb{P}^2$  elements for the velocity and  $\mathbb{P}^1$  elements for the pressure.

# Domain decomposition

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**MPI** domain decomposition with graph partitioner. **PETSc** matrices and solvers.

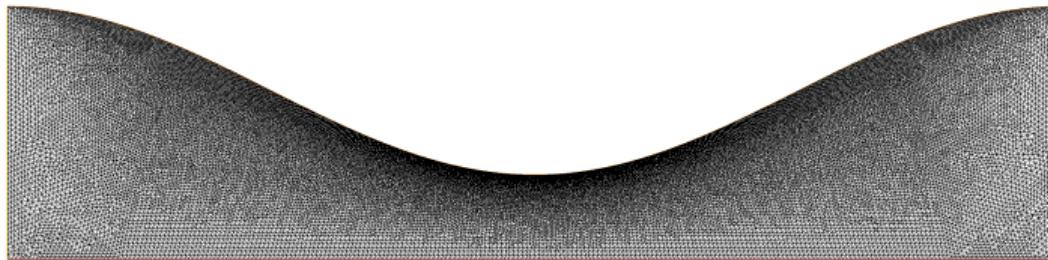
Example with **6 domains**:



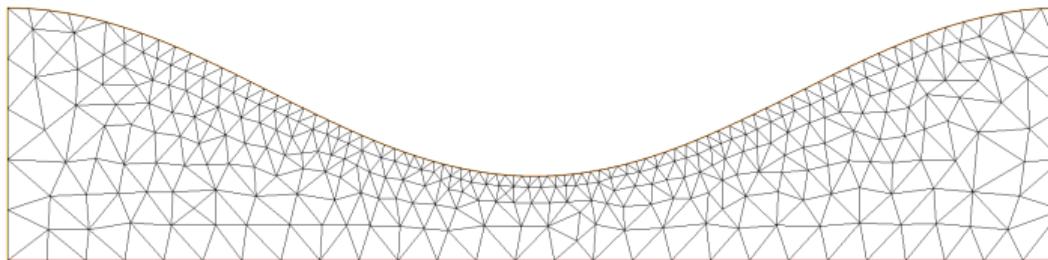
# Geometric multigrid

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We use a **geometric multigrid** solver for fast convergence using a large number of MPI processes.



Level 0 (fine)



Level 2 (coarse)

## Timings and #dofs

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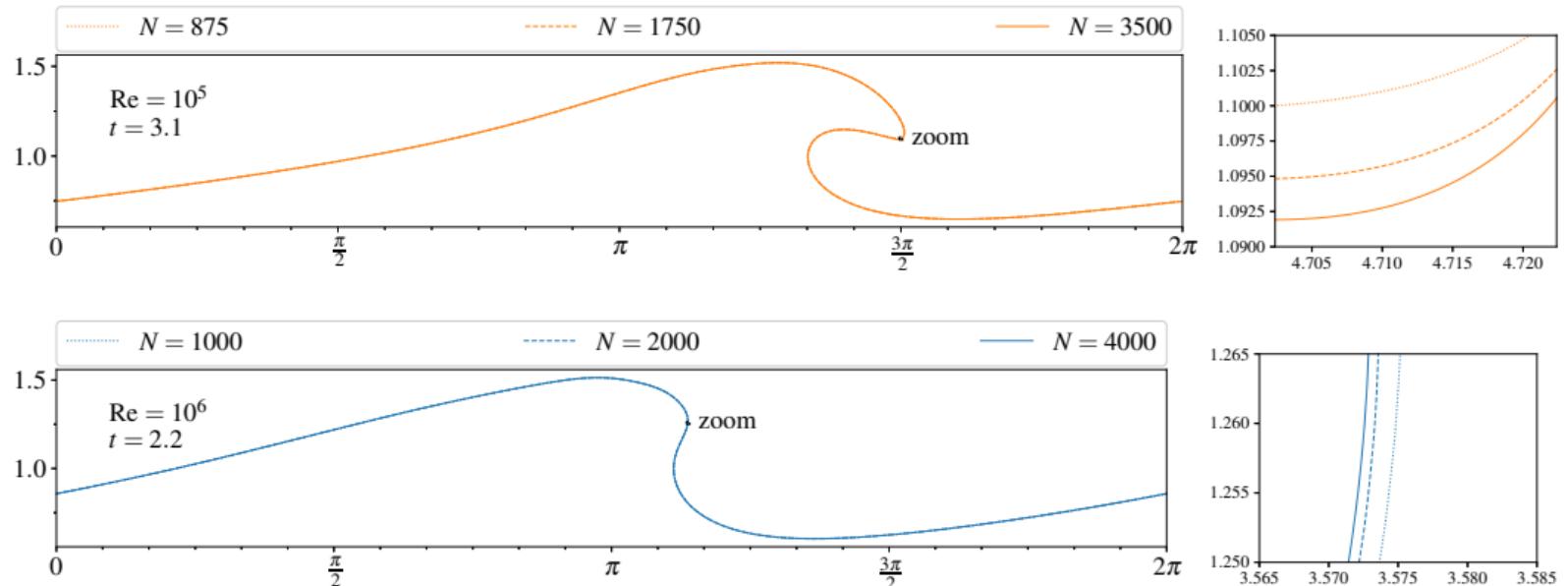
Between 1 and 3.5 million unknowns, convergence in  $\sim 5 \pm 2$  GMRES iterations to machine precision  $\rightarrow \sim 20$  seconds on 48 CPU cores.

Main computational limitations due to **FreeFem memory management** and **spurious behaviors in mesh handling**.

Clearly this **method is not efficient enough to handle 3d!**

Re	N	# dofs at the start			# dofs at the end		
		# triangles	( $u, p$ )	w	# triangles	( $u, p$ )	w
$10^2$	3000	195,314	886,913	198,514	780,572	3,520,574	783,722
$10^3$	3500	221,748	1,007,116	225,448	615,382	2,675,923	596,294
$10^4$	3500	221,368	1,005,406	225,068	510,266	2,305,447	513,966
$10^5$	3500	222,970	1,012,615	226,670	450,648	2,037,166	454,348
$10^6$	4000	272,948	1,238,766	277,148	498,250	2,252,625	502,450

# Convergence



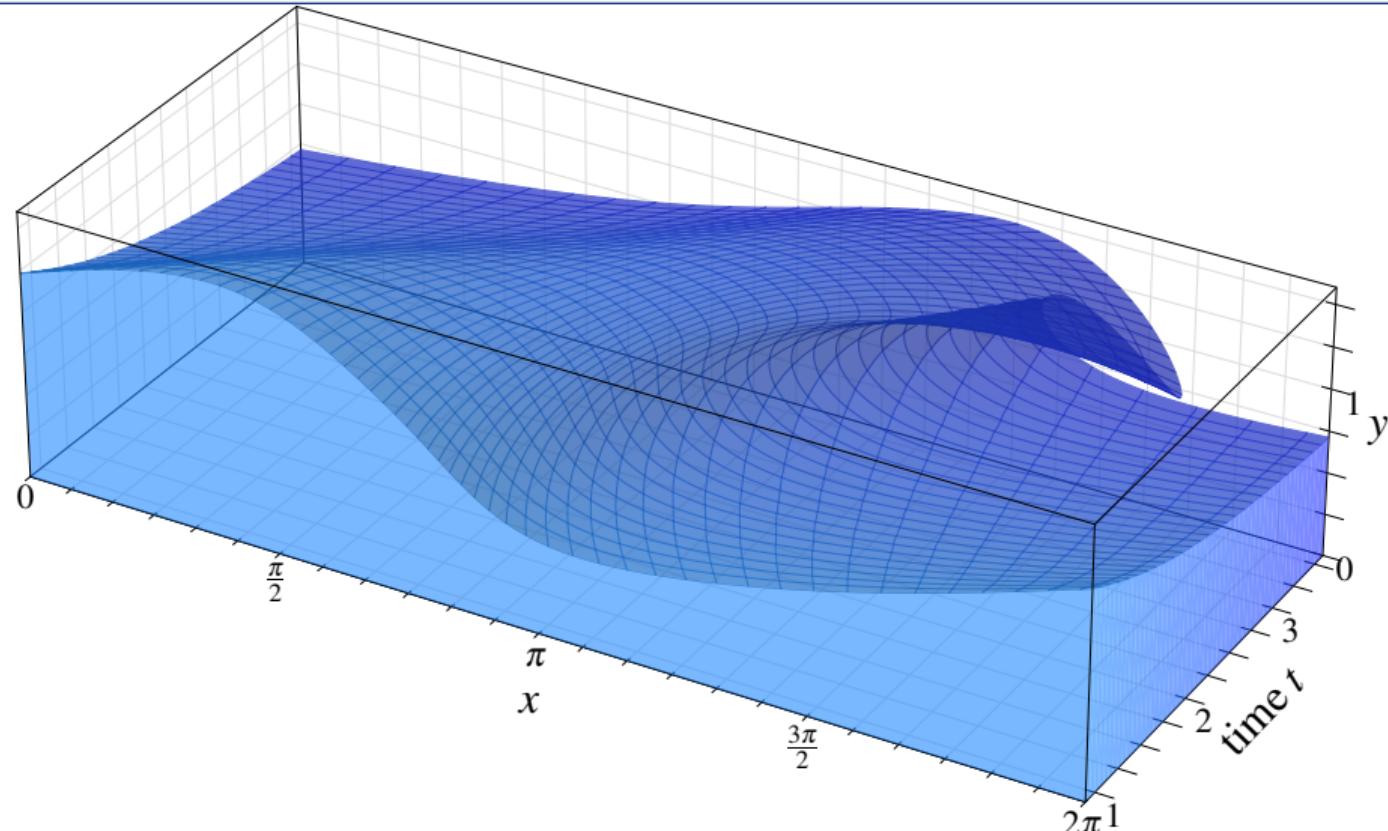
## The $\text{Re} \rightarrow +\infty$ limit

Comparing our results with the inviscid solution



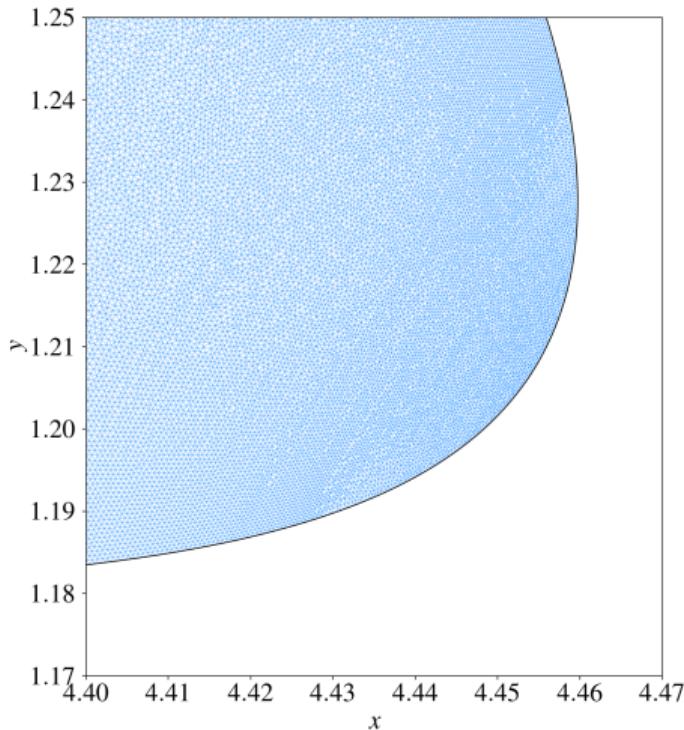
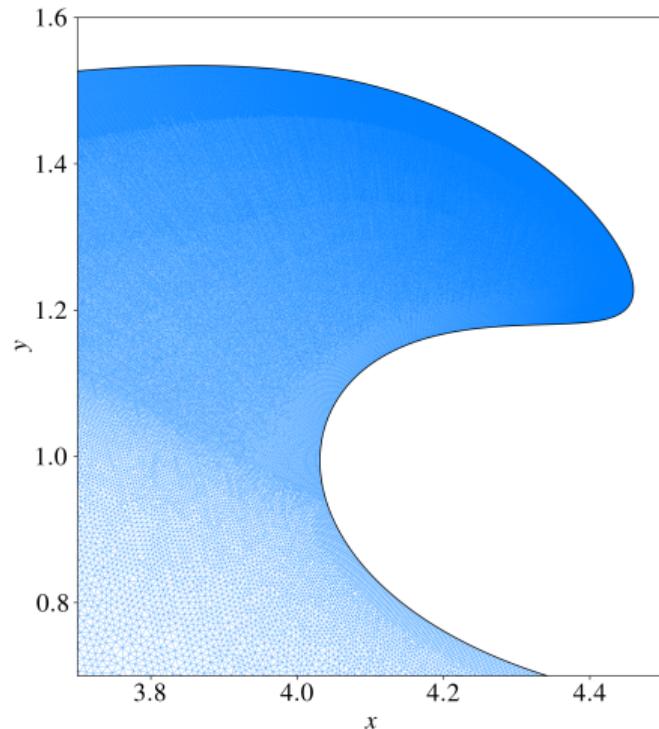
# $\text{Re} = 10^6$ **result**

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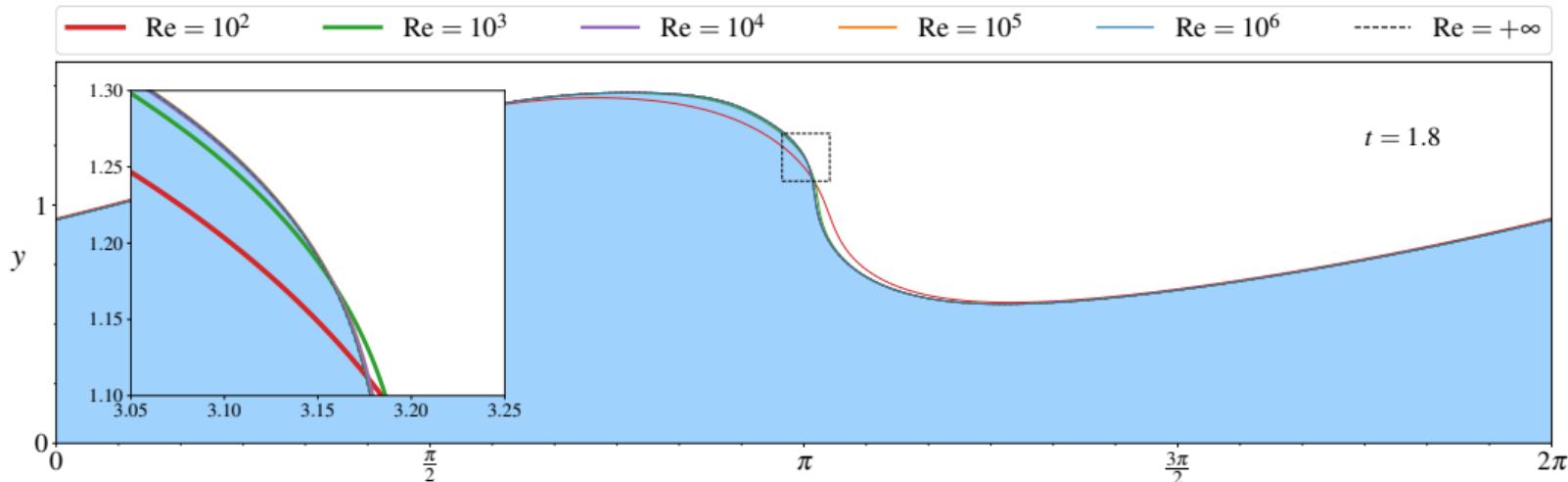


# Mesh at $\text{Re} = 10^6$

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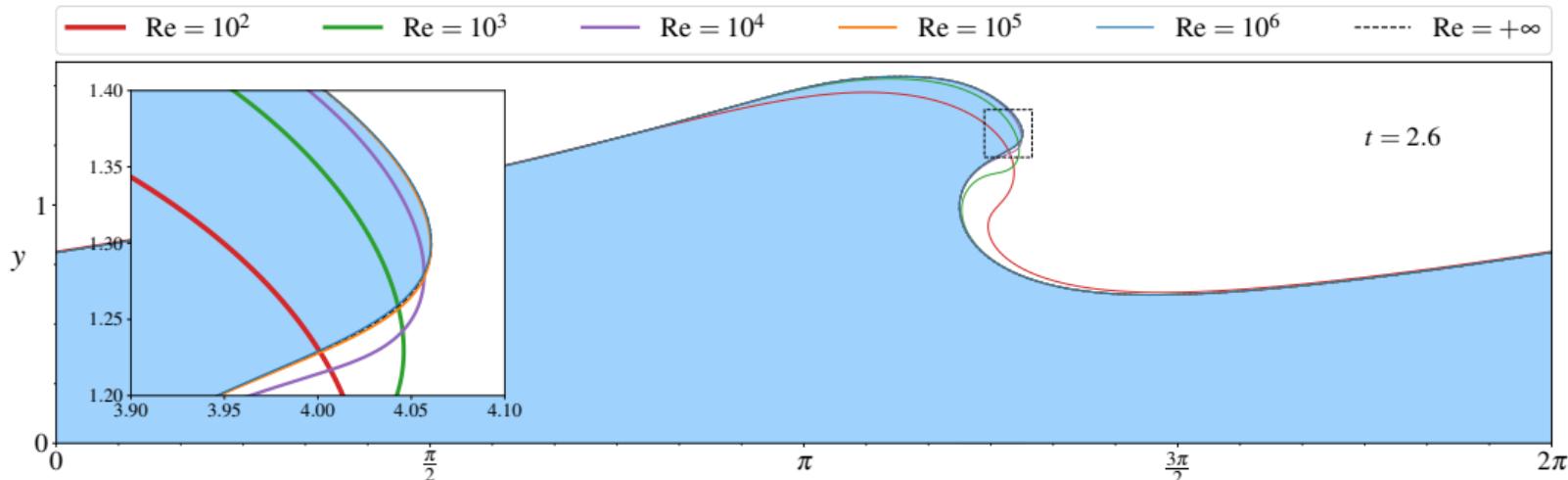


# Interface for different values of Re



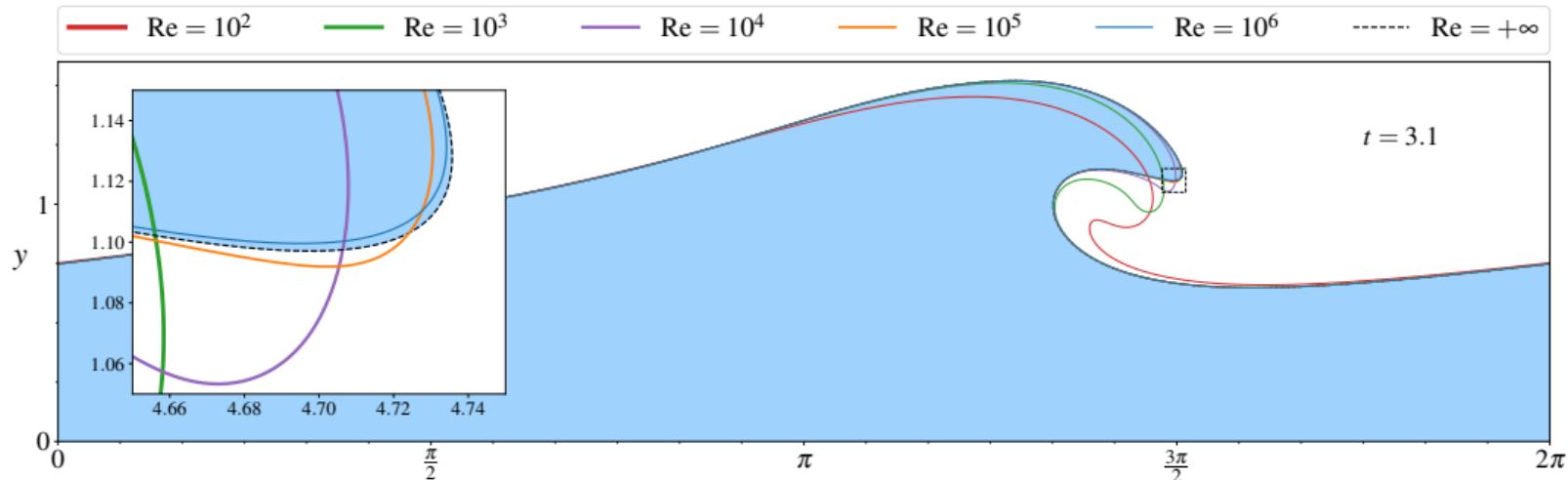
Re =  $+\infty$  simulations (i.e. Euler solution) computed with the numerical methods of [Dormy & Lacave \(2024\)](#).

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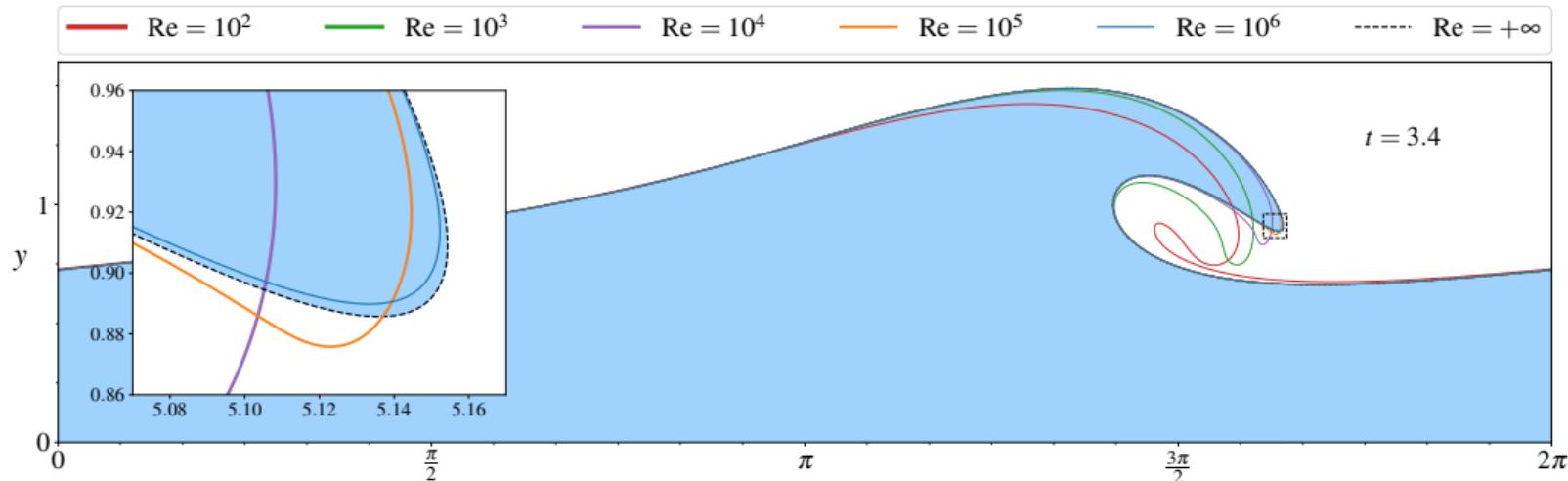
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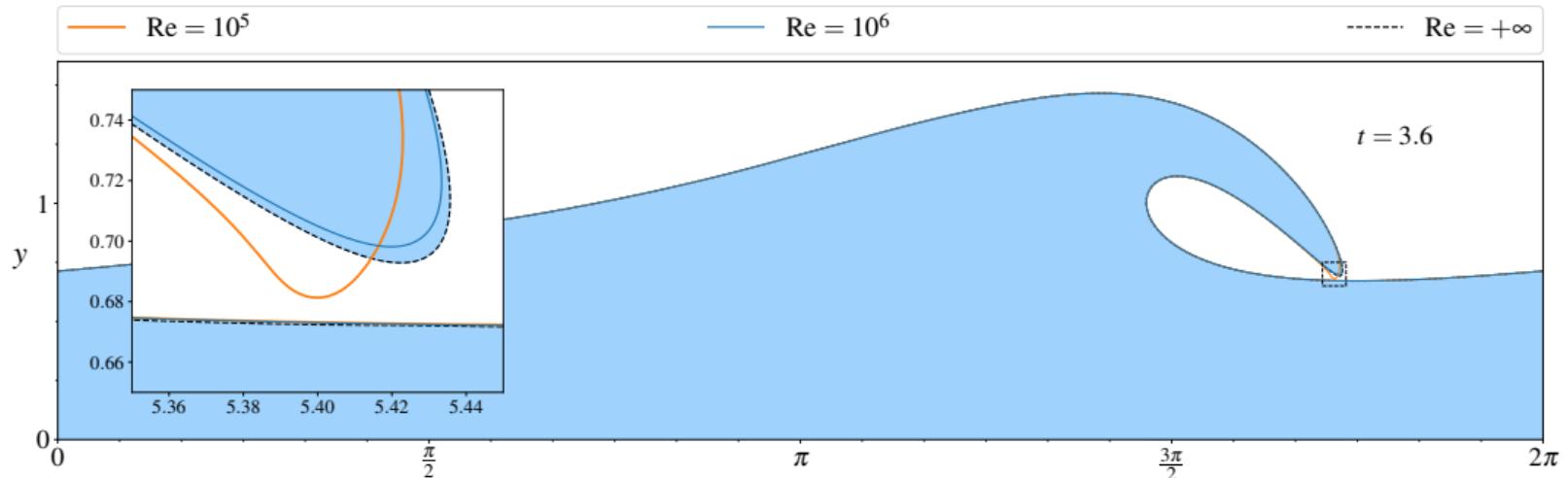
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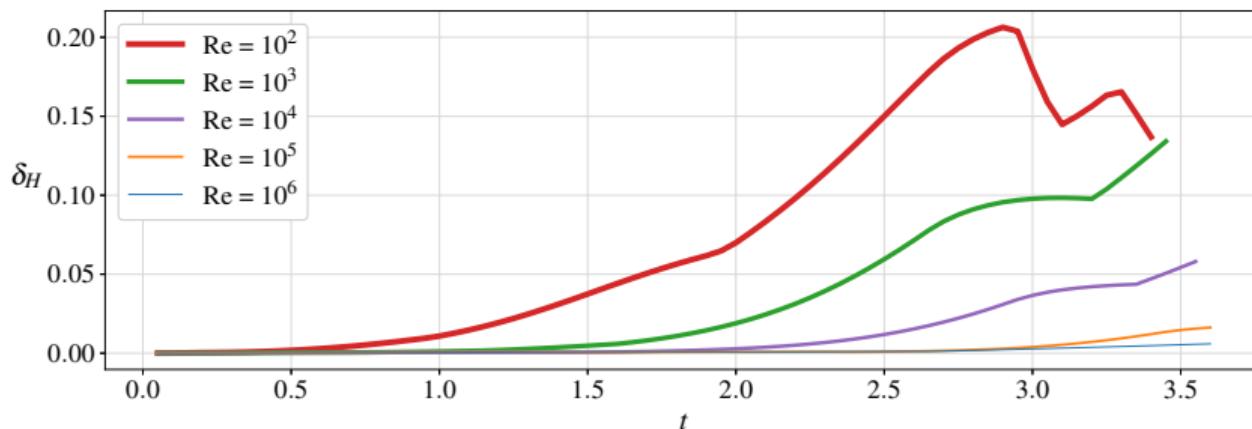
# Convergence

Haussdorff distance

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To compare the Euler solution with Navier-Stokes, we use the **Haussdorff distance**,

$$\delta_H(\gamma_1, \gamma_2) = \max \left\{ \tilde{\delta}_H(\gamma_1, \gamma_2), \tilde{\delta}_H(\gamma_2, \gamma_1) \right\} \quad \text{where} \quad \tilde{\delta}_H(\gamma_1, \gamma_2) = \max_{s_1} \min_{s_2} |\gamma_1(s_1) - \gamma_2(s_2)|$$



# Boundary Layer

Where does the viscous dissipation happen?



# Energy considerations

Link with vorticity

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The **kinetic energy equation** can be obtained multiplying Navier-Stokes eq. by  $\mathbf{u}$ ,

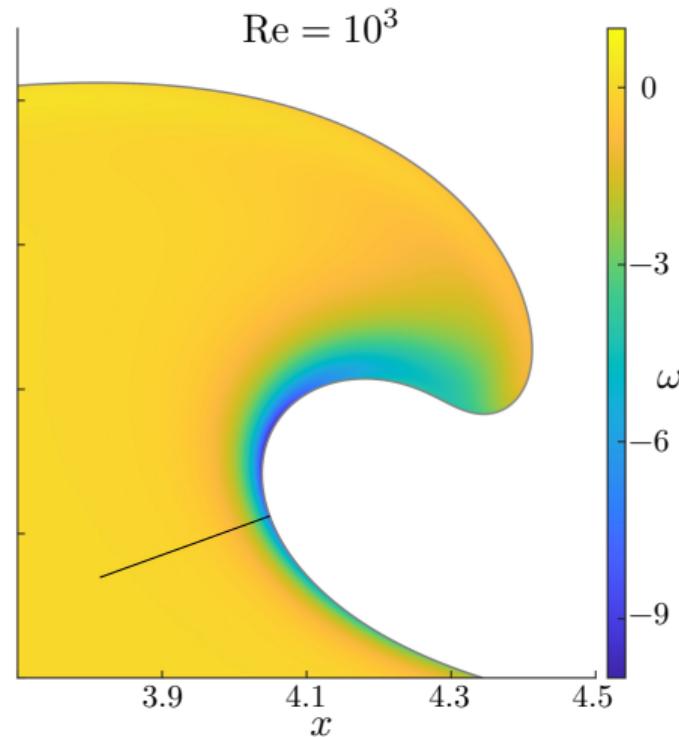
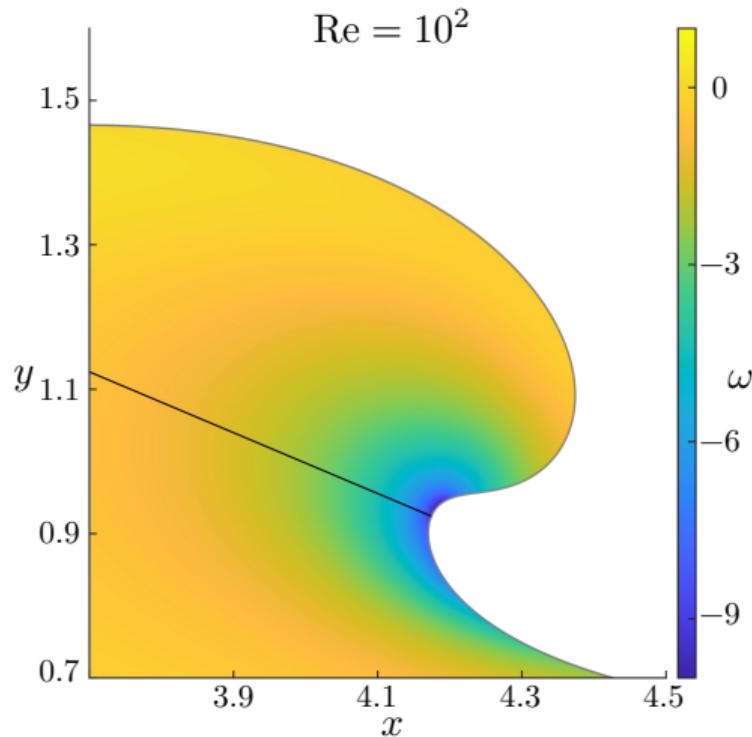
$$\partial_t \left( \frac{\mathbf{u}^2}{2} \right) = \mathbf{g} \cdot \mathbf{u} - \mathbf{u} \cdot \nabla p + \frac{1}{\text{Re}} [\nabla \cdot (\mathbf{u}^\perp \boldsymbol{\omega}) - \boldsymbol{\omega}^2]$$

where  $\mathbf{u}^\perp = [-u_y, u_x]$  and  $\boldsymbol{\omega} = \nabla^\perp \cdot \mathbf{u}$  is the **vorticity**.

This shows that **fluids dissipates energy in the support of  $\boldsymbol{\omega}$ !**

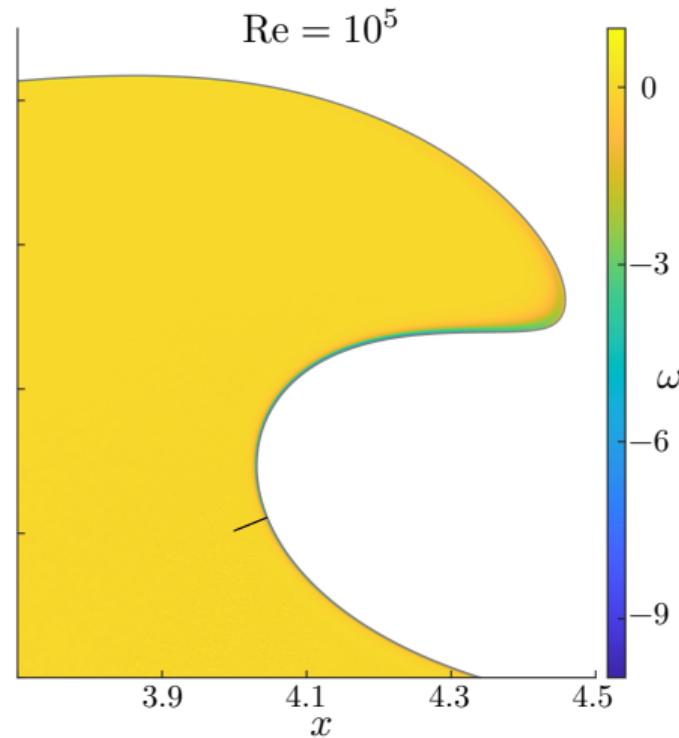
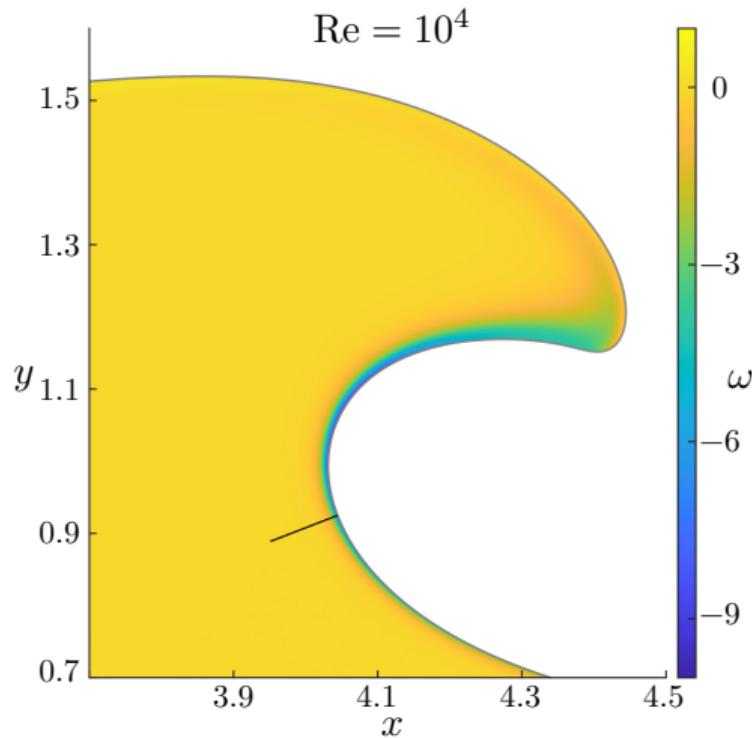
# Viscous dissipation

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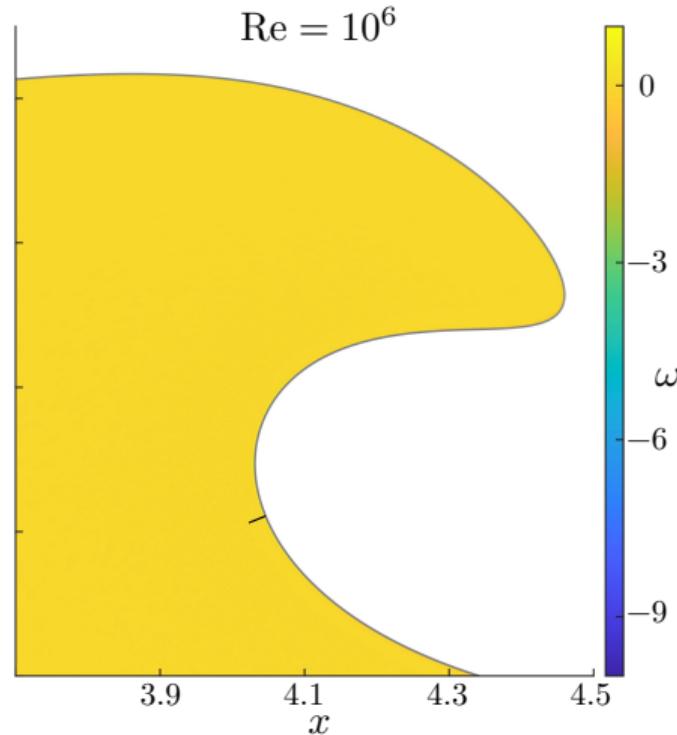
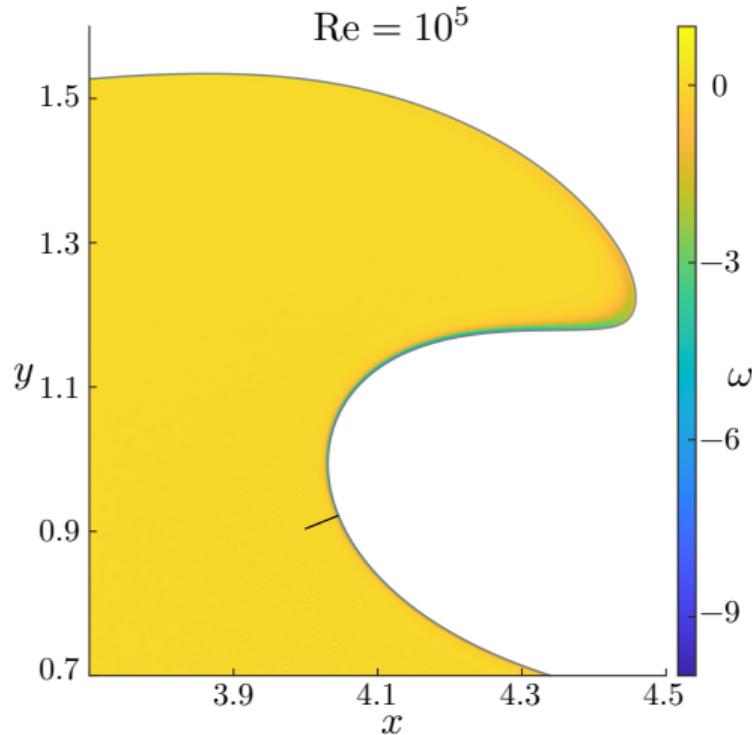
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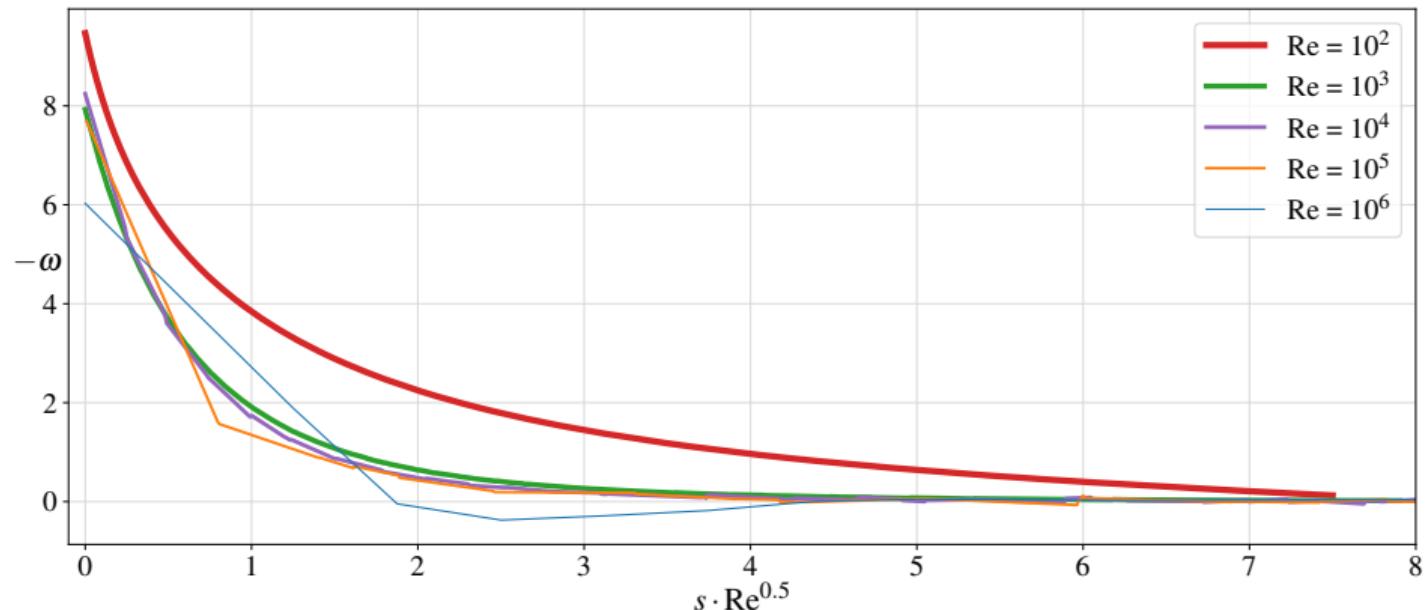
# Viscous dissipation

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# Size of the boundary layer

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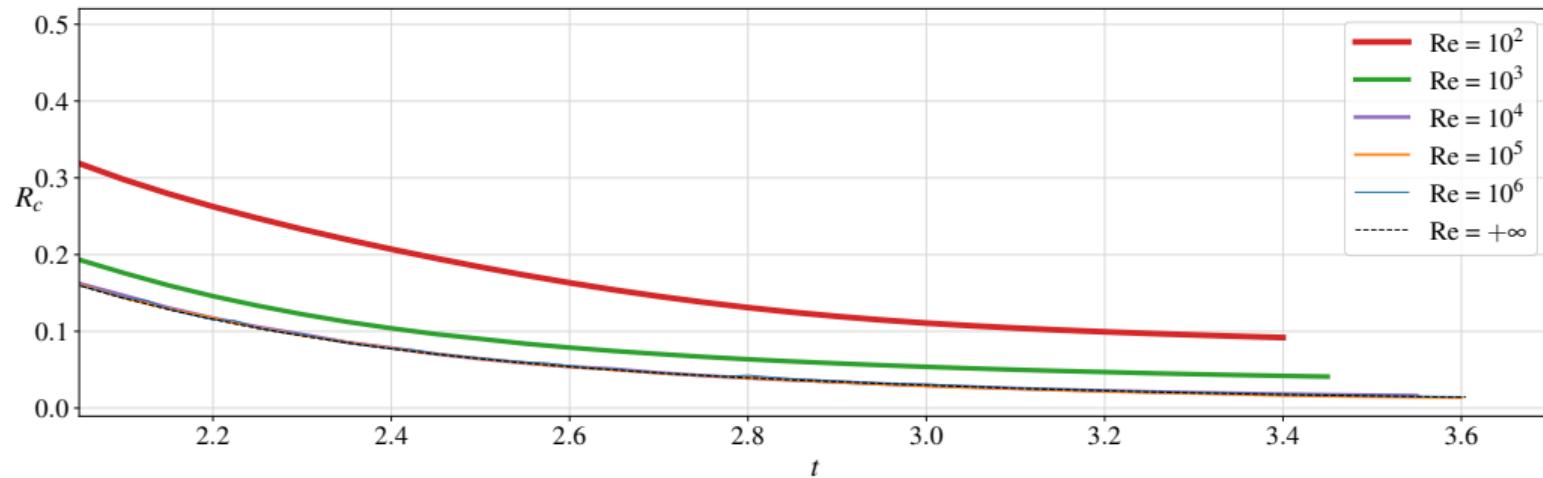


Exhibits a  $Re^{-\frac{1}{2}}$  scaling (usual in BL theory).

Thank you!

# Maximum curvature of the interface

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where  $R_C = \kappa^{-1}$  is the curvature radius.